

(Probabilités)

Calculus on Fock space and a non-adapted quantum Itô formula

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Abstract - *The aim of this note is to introduce a calculus on Fock space with its probabilistic interpretations, and to give a detailed presentation of the associated quantum Itô formula.*

Calcul sur l'espace de Fock et une formule d'Itô non-commutative anticipante

Résumé - Le but de cette note est d'introduire un calcul sur l'espace de Fock et de donner une présentation détaillée de la formule d'Itô non-commutative associée.

Version française abrégée - Soit $\Phi = \bigoplus_{n \geq 0} H^{\otimes n}$ l'espace de Fock symétrique sur $H = L^2(\mathbb{R}_+)$ avec ses opérateurs ∇, ∇^* d'annihilation et de création. Soit \mathcal{S} le sous-ensemble de Φ dont les éléments ont un développement fini qui ne fait intervenir que des fonctions \mathcal{C}_c^1 , et soit \mathcal{U} l'ensemble des éléments de $\Phi \otimes L^2(\mathbb{R}_+)$ de la forme $\sum_{i=1}^n F_i \otimes h_i$, avec $h_1, \dots, h_n \in \mathcal{C}_c^1(\mathbb{R}_+)$, et $F_1, \dots, F_n \in \mathcal{S}$, $n \in \mathbb{N}$. Pour $h \in L^2(\mathbb{R}_+)$, on pose $\overset{\circ}{h}(t) = \int_0^t h(s) ds$ et $h|_t = 1_{[t, \infty[} h$, $t \in \mathbb{R}_+$.

Définition On définit les opérateurs linéaires $\tilde{\nabla} : \Phi \rightarrow \Phi \otimes L^2(\mathbb{R}_+)$ sur \mathcal{S} et $\tilde{\nabla}^* : \Phi \otimes L^2(\mathbb{R}_+) \rightarrow \Phi$ sur \mathcal{U} par

$$\begin{aligned} \tilde{\nabla}_t f^{\otimes n} &= -n f|_t' \circ f^{\otimes(n-1)}, \quad t \in \mathbb{R}_+, f \in \mathcal{C}_c^1(\mathbb{R}_+), \quad n \in \mathbb{N}, \\ \tilde{\nabla}^*(f^{\otimes n} \otimes h) &= n \left(f \overset{\circ}{h} \right)' \circ f^{\otimes(n-1)}, \quad f, h \in \mathcal{C}_c^1(\mathbb{R}_+), \end{aligned}$$

et par polarisation.

Les opérateurs $\tilde{\nabla}, \tilde{\nabla}^*$ sont fermables et $\tilde{\nabla}^* : \Phi \otimes L^2(\mathbb{R}_+) \rightarrow \Phi$ est l'adjoint de $\tilde{\nabla} : \Phi \rightarrow \Phi \otimes L^2(\mathbb{R}_+)$. Soit $(W, L^2(\mathbb{R}_+), \mu)$ l'espace de Wiener avec le mouvement brownien $(B_t)_{t \in \mathbb{R}_+}$. Pour $h \in L^2(\mathbb{R}_+)$ avec $\sup_{x \in \mathbb{R}_+} |h(x)| < 1$, soit $\nu_h(t) = t + \int_0^t h(s) ds$, $t \in \mathbb{R}_+$. On définit le changement de temps $\mathcal{T}_h : W \rightarrow W$ par (1). Sous l'identification de Wiener usuelle entre Φ et $L^2(W, \mu)$,

$$\int_0^\infty h(t) \left(\tilde{\nabla}_t + \frac{1}{2} \nabla_t \nabla_t \right) F dt = - \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (F \circ \mathcal{T}_{\varepsilon h} - F), \quad F \in \mathcal{S}.$$

Dans l'interprétation poissonnienne de Φ , $\nabla + \tilde{\nabla}$ coïncide avec l'opérateur \tilde{D} défini sur l'espace de Poisson par changement de temps (4). Soit $\tilde{\delta}$ l'adjoint de \tilde{D} , et soient $(\tilde{a}_t^\pm)_{t \in \mathbb{R}_+}$ les processus d'opérateurs définis par

$$\tilde{a}_t^- F = \int_0^t \tilde{D}_s F ds, \quad \tilde{a}_t^+ F = \tilde{\delta}(1_{[0,t]} F), \quad F \in \mathcal{S}, \quad t \in \mathbb{R}_+.$$

Comme le montre la définition de $\tilde{\nabla}$, $(\tilde{a}_t^\pm)_{t \in \mathbb{R}_+}$ sont des perturbations des processus de création et d'annihilation sur Φ , et ne sont pas adaptés en tant que processus d'opérateurs. Par conséquent \tilde{a}_t^\pm ne commute pas avec sa différentielle, et la formule d'Itô de [1] doit donc être écrite comme

Proposition Si X, Y sont des processus simples adaptés tels que $\Xi \subset \text{Dom}(X_s Y_s)$, $s \in \mathbb{R}_+$:

$$\begin{aligned} \int_0^t X_s d\tilde{a}_s^\varepsilon \int_0^t Y_s d\tilde{a}_s^\eta &= \int_0^t d\tilde{a}_s^\varepsilon X_s \left(\int_0^s Y_u d\tilde{a}_u^\eta \right) + \int_0^t \left(\int_0^s X_u d\tilde{a}_u^\varepsilon \right) Y_s d\tilde{a}_s^\eta \\ &\quad + \int_0^t X_s Y_s d\tilde{a}_s^\varepsilon \cdot d\tilde{a}_s^\eta, \quad \varepsilon, \eta = -, +, \end{aligned}$$

où le produit $d\tilde{a}_s^- \cdot d\tilde{a}_s^+$ vaut dN_t , les autres produits étant nuls.

1 A calculus on Fock space

The annihilation operator $\nabla : \Phi \rightarrow \Phi \otimes H$ on the symmetric Fock space $\Phi = \bigoplus_{n \geq 0} H^{\circ n}$ over $H = L^2(\mathbb{R}_+)$ is defined by transformation of the tensor $f^{\circ n}$, $f \in H$, into $\nabla f^{\circ n} = n f^{\circ(n-1)} \otimes f$, $n \in \mathbb{N}$, while the creation operator $\nabla^* : \Phi \otimes H \rightarrow \Phi$ satisfies $\nabla^* f^{\circ n} \otimes g = f^{\circ n} \circ g$, $n \in \mathbb{N}$. The space Φ has two main probabilistic interpretations. In the Wiener interpretation, the sum of the creation and annihilation operators is identified to a Brownian motion which can be perturbed by the number operator process in order to give a Poisson process in the Poisson interpretation of Φ . The operators $\tilde{\nabla}$ and $\tilde{\nabla}^*$ defined below will be interpreted as perturbations of ∇ and ∇^* . For $f \in \mathcal{C}_c^1(\mathbb{R}_+)$, $h \in L^2(\mathbb{R}_+)$, let $\overset{\circ}{h}(t) = \int_0^t h(s) ds$, $h|_t = 1_{[t, \infty[} h$, and $f'(t) = \frac{d}{dt} f(t)$, $t \in \mathbb{R}_+$. Let \mathcal{S} denote the subset of Φ whose elements have a finite development involving only \mathcal{C}_c^1 functions, and let \mathcal{U} denote the set of elements of $\Phi \otimes L^2(\mathbb{R}_+)$ of the form $\sum_{i=1}^n F_i \otimes h_i$, with $h_1, \dots, h_n \in \mathcal{C}_c^1(\mathbb{R}_+)$, and $F_1, \dots, F_n \in \mathcal{S}$, $n \in \mathbb{N}$.

Definition 1 We define the linear operators $\tilde{\nabla} : \Phi \rightarrow \Phi \otimes L^2(\mathbb{R}_+)$ on \mathcal{S} and $\tilde{\nabla}^* : \Phi \otimes L^2(\mathbb{R}_+) \rightarrow \Phi$ on \mathcal{U} by

$$\begin{aligned} \tilde{\nabla}_t f^{\circ n} &= -n f'_t \circ f^{\circ(n-1)}, \quad t \in \mathbb{R}_+, f \in \mathcal{C}_c^1(\mathbb{R}_+), \quad n \in \mathbb{N}, \\ \tilde{\nabla}^*(f^{\circ n} \otimes h) &= n \left(f \overset{\circ}{h} \right)' \circ f^{\circ(n-1)}, \quad f, h \in \mathcal{C}_c^1(\mathbb{R}_+), \end{aligned}$$

and by polarization of these expressions.

Both $\tilde{\nabla}$, $\tilde{\nabla}^*$ are closable, and the operator $\tilde{\nabla}^* : \Phi \otimes L^2(\mathbb{R}_+) \rightarrow \Phi$ is the adjoint operator of $\tilde{\nabla} : \Phi \rightarrow \Phi \otimes L^2(\mathbb{R}_+)$. The domain of $\tilde{\nabla}$ is determined by the equality of norms

$$\| \tilde{\nabla} F \|_{\Phi \otimes L^2(\mathbb{R}_+)}^2 = \sum_{n \geq 1} n^3 \int_0^\infty t \| \partial_1 f_n(t, \cdot) \|_{L^2(\mathbb{R}_+)^{\circ(n-1)}}^2 dt, \quad F \in \mathcal{S}, F = \sum_{n \in \mathbb{N}} f_n.$$

2 Probabilistic interpretations

Let $(W, L^2(\mathbb{R}_+), \mu)$ denote the classical Wiener space, with Brownian motion $(B_t)_{t \in \mathbb{R}_+}$. We recall that under the usual identification between Φ and $L^2(W, \mu)$, ∇ is identified to a derivation operator which satisfies

$$(\nabla F, h)_{L^2(\mathbb{R}_+)} = \lim_{\varepsilon \rightarrow 0} \frac{F(B. + \varepsilon \int_0^\cdot h(s) ds) - F}{\varepsilon}, \quad F \in \mathcal{S}, h \in L^2(\mathbb{R}_+),$$

cf. e.g. [2], [3], [4]. For $h \in L^2(\mathbb{R}_+)$, with $\sup_{x \in \mathbb{R}_+} |h(x)| < 1$, let $\nu_h(t) = t + \int_0^t h(s) ds$, $t \in \mathbb{R}_+$. We define a mapping $\mathcal{T}_h : W \rightarrow W$, $t, \varepsilon \in \mathbb{R}_+$, as

$$\mathcal{T}_h(\omega) = \omega \circ \nu_h^{-1}, \quad h \in L^2(\mathbb{R}_+), \sup_{x \in \mathbb{R}_+} |h(x)| < 1. \quad (1)$$

Although \mathcal{T}_h is not absolutely continuous on the Wiener space, the functional $F \circ \mathcal{T}_h$ is well-defined for $F \in \mathcal{S}$, since elements of \mathcal{S} can be defined trajectory by trajectory.

Proposition 1 *On the Wiener space, $\tilde{\nabla}$ satisfies the relation*

$$\tilde{\nabla}_t(FG) = F\tilde{\nabla}_tG + G\tilde{\nabla}_tF - \nabla_tF\nabla_tG, \quad t \in \mathbf{R}_+, \quad F, G \in \mathcal{S}, \quad (2)$$

and

$$\int_0^\infty h(t) \left(\tilde{\nabla}_t + \frac{1}{2} \nabla_t \nabla_t \right) F dt = - \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (F \circ \mathcal{T}_{\varepsilon h} - F), \quad F \in \mathcal{S}, \quad h \in L^2(\mathbf{R}_+).$$

Let $(N_t)_{t \in \mathbf{R}_+}$ be a Poisson process on a probability space (B, P) . From [5], [6], [7], ∇_t acts in the Poisson identification of $L^2(B, P)$ with Φ by perturbation of the Poisson process trajectory via addition of a jump at time t , and

$$\nabla(FG) = F\nabla G + G\nabla F + \nabla F\nabla G, \quad F, G \in \mathcal{S}. \quad (3)$$

A gradient operator $\tilde{D} : L^2(B) \rightarrow L^2(B) \otimes L^2(\mathbf{R}_+)$ on Poisson space, cf. [8], [9], is defined as

$$(\tilde{D}F, h)_{L^2(\mathbf{R}_+)} = - \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (F \circ \mathcal{T}_{\varepsilon h} - F), \quad F \in \mathcal{S}, \quad h \in L^2(\mathbf{R}_+), \quad (4)$$

where the transformation \mathcal{T}_h is defined by the time change ν_h applied to $(N_t)_{t \in \mathbf{R}_+}$. Unlike in the Wiener space case, the transformation \mathcal{T}_h is absolutely continuous in the Poisson space case. The operator \tilde{D} is closable, its adjoint $\tilde{\delta}$ satisfies

$$\tilde{\delta}(u) = \int_0^\infty u(s) d(N_s - s) - \int_0^\infty \tilde{D}_s u(s) ds,$$

for $u(t) = f(t, T_1, \dots, T_n)$, $f \in \mathcal{C}_c^1(\mathbf{R}_+^{n+1})$, and coincides with the compensated Poisson stochastic integral on the square-integrable adapted processes, cf. [8], [9].

Proposition 2 *On the Poisson space, $\tilde{D} = \nabla + \tilde{\nabla}$, and $\tilde{\nabla}$ also satisfies the relation*

$$\tilde{\nabla}_t(FG) = F\tilde{\nabla}_tG + G\tilde{\nabla}_tF - \nabla_tF\nabla_tG, \quad t \in \mathbf{R}_+, \quad F, G \in \mathcal{S}.$$

In the Wiener interpretation of Φ , $\tilde{\delta}$ is identified to an extension $\nabla^* + \tilde{\nabla}^*$ of the stochastic integral with respect to $(B_t)_{t \in \mathbf{R}_+}$. As an application of this calculus, we obtain the following absolute continuity criterion for single Poisson stochastic integrals.

Proposition 3 *Let $f \in L^2(\mathbf{R}_+)$ such that $\int_0^\infty t f'(t)^2 dt < \infty$ and $\{f' = 0\}$ has finite Lebesgue measure. Then the law of $\int_0^\infty f(t) d(N_t - t)$ is absolutely continuous with respect to the Lebesgue measure.*

3 A non-adapted quantum Itô formula

Define the operator processes $(\tilde{a}_t^\pm)_{t \in \mathbf{R}_+}$ as

$$\tilde{a}_t^- F = \int_0^t \tilde{D}_s F ds, \quad \tilde{a}_t^+ F = \tilde{\delta}(1_{[0,t]} F), \quad F \in \mathcal{S}, \quad t \in \mathbf{R}_+.$$

In the Poisson interpretation of Φ , $(\tilde{a}_t^- + \tilde{a}_t^+)_{t \in \mathbf{R}_+}$ is a decomposition of the multiplication operator by the compensated Poisson process $(N_t - t)_{t \in \mathbf{R}_+}$. The aim of this section is to

make a precise statement of a fact that has been overlooked in [1], [10]. Namely, the processes $(\tilde{a}_t^\pm)_{t \in \mathbf{R}_+}$ are not adapted operator processes, as can be seen from the definition of $\tilde{\nabla}$. Consequently, \tilde{a}_t^\pm does not commute with its differential. Let

$$\int_0^t d\tilde{a}_s^+ \tilde{a}_s^\pm F = \tilde{\delta} \left(1_{[0,t]} \tilde{a}_s^\pm F \right), \quad \int_0^t d\tilde{a}_s^- \tilde{a}_s^\pm F = \int_0^t \tilde{D}_s \tilde{a}_s^\pm F ds, \quad F \in \mathcal{S}.$$

The integrals $\int_0^t \tilde{a}_s^\pm d\tilde{a}_s^\pm$ and $\int_0^t \tilde{a}_s^\pm d\tilde{a}_s^\mp$ are defined by duality on \mathcal{S} as being distribution-valued in the dual $\mathbb{D}_{2,-1}$ of $\text{Dom}(\tilde{D})$. They are the respective adjoints of $\int_0^t d\tilde{a}_s^- \tilde{a}_s^\mp$, $\int_0^t d\tilde{a}_s^+ \tilde{a}_s^\mp$. Prop. 7 in [1] should read after reordering of the differentials:

Proposition 4 *We have for simple adapted processes X, Y such that $\Xi \subset \text{Dom}(X_s Y_s)$, $s \in \mathbf{R}_+$:*

$$\begin{aligned} \int_0^t X_s d\tilde{a}_s^\varepsilon \int_0^t Y_s d\tilde{a}_s^\eta &= \int_0^t d\tilde{a}_s^\varepsilon X_s \left(\int_0^s Y_u d\tilde{a}_u^\eta \right) + \int_0^t \left(\int_0^s X_u d\tilde{a}_u^\varepsilon \right) Y_s d\tilde{a}_s^\eta \\ &\quad + \int_0^t X_s Y_s d\tilde{a}_s^\varepsilon \cdot d\tilde{a}_s^\eta, \quad \varepsilon, \eta = -, +, \end{aligned}$$

where the product $d\tilde{a}_s^- \cdot d\tilde{a}_s^+$ is given by $d\tilde{a}_t^- \cdot d\tilde{a}_t^+ = dN_t$, the other products being zero.

This expression might not always be defined on all $F \in \mathcal{S}$, and we will prove in general the following relations:

$$\begin{aligned} &\langle \int_0^t Y_s d\tilde{a}_s^+ G, \int_0^t X_s^* d\tilde{a}_s^- F \rangle \\ &= \int_0^t \langle \int_0^s Y_u d\tilde{a}_u^+ G, X_s^* \tilde{D}_s F \rangle ds + \int_0^t \langle Y_s G, \tilde{D}_s \int_0^s X_u^* d\tilde{a}_u^- F \rangle ds, \\ &\langle \int_0^t Y_s d\tilde{a}_s^- G, \int_0^t X_s^* d\tilde{a}_s^+ F \rangle \\ &= \int_0^t \langle \tilde{D}_s G, \int_0^s Y_u^* d\tilde{a}_u^+ X_s^* F \rangle ds + \int_0^t \langle \tilde{D}_s \int_0^s X_u d\tilde{a}_u^- Y_s G, F \rangle ds, \\ &\langle \int_0^t Y_s d\tilde{a}_s^+ G, \int_0^t X_s^* d\tilde{a}_s^+ F \rangle \\ &= \int_0^t \langle \tilde{D}_s X_s \int_0^s Y_u d\tilde{a}_u^+ G, F \rangle ds + \int_0^t \langle Y_s G, \int_0^s X_u^* d\tilde{a}_u^+ \tilde{D}_s F \rangle ds, \\ &\quad + \langle \int_0^t X_s Y_s G dN_s, F \rangle, \\ &\langle \int_0^t Y_s d\tilde{a}_s^- G, \int_0^t X_s^* d\tilde{a}_s^- F \rangle \\ &= \int_0^t \langle \int_0^s Y_u d\tilde{a}_u^- G, X_s^* \tilde{D}_s F \rangle ds + \int_0^t \langle Y_s \tilde{D}_s G, \int_0^s X_u^* d\tilde{a}_u^- F \rangle ds, \end{aligned}$$

for $F, G \in \mathcal{S}$. By non-commutative linearity and adaptedness of $(X_t)_{t \in \mathbf{R}_+}$, $(Y_t)_{t \in \mathbf{R}_+}$, they are a consequence of the following lemma, which is identical to Lemma 1 of [1], except for the reordering between \tilde{a}_s^\pm and its differential.

Lemma 1 *We have*

$$\tilde{a}_t^+ \tilde{a}_t^- = \int_0^t d\tilde{a}_s^+ \tilde{a}_s^- + \int_0^t \tilde{a}_s^+ d\tilde{a}_s^-, \quad (5)$$

$$\tilde{a}_t^\pm \tilde{a}_t^\pm = \int_0^t \tilde{a}_s^\pm d\tilde{a}_s^\pm + \int_0^t d\tilde{a}_s^\pm \tilde{a}_s^\pm, \quad (6)$$

$$\tilde{a}_t^- \tilde{a}_t^+ = \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ + \int_0^t \tilde{a}_s^- d\tilde{a}_s^+ + N_t. \quad (7)$$

Proof. We have for (5):

$$\begin{aligned}
\langle \tilde{a}_t^- G, \tilde{a}_t^- F \rangle &= \langle \int_0^t \tilde{D}_u G du, \int_0^t \tilde{D}_s F ds \rangle \\
&= \int_0^t \int_0^s \langle \tilde{D}_u G, \tilde{D}_s F \rangle dud s + \int_0^t \int_0^u \langle \tilde{D}_u G, \tilde{D}_s F \rangle dud s \\
&= \langle \tilde{\delta} \left(1_{[0,t]}(\cdot) \int_0^\cdot \tilde{D}_u G du \right), F \rangle + \langle G, \tilde{\delta} \left(1_{[0,t]}(\cdot) \int_0^\cdot \tilde{D}_u F du \right) \rangle \\
&= \langle \int_0^t d\tilde{a}_s^+ \tilde{a}_s^- G, F \rangle + \langle G, \int_0^t d\tilde{a}_s^+ \tilde{a}_s^- F \rangle \\
&= \langle G, \left(\int_0^t \tilde{a}_s^+ d\tilde{a}_s^- + \int_0^t d\tilde{a}_s^+ \tilde{a}_s^- \right) F \rangle, \quad F, G \in \mathcal{S},
\end{aligned}$$

and

$$\begin{aligned}
\langle \tilde{a}_t^+ F, \tilde{a}_t^- G \rangle &= \int_0^t \langle \tilde{a}_s^+ F, \tilde{D}_s G \rangle ds + \langle F, \int_0^t \tilde{a}_s^- G d\tilde{N}_s \rangle - \int_0^t \langle \tilde{D}_s F, \tilde{a}_s^- G \rangle ds \\
&= \int_0^t \langle \tilde{a}_s^+ F, \tilde{D}_s G \rangle ds + \langle F, \int_0^t d\tilde{a}_s^+ \tilde{a}_s^- G \rangle + \langle F, \int_0^t d\tilde{a}_s^- \tilde{a}_s^- G \rangle - \int_0^t \langle \tilde{D}_s F, \tilde{a}_s^- G \rangle ds \\
&= \int_0^t \langle \tilde{a}_s^+ F, \tilde{D}_s G \rangle ds + \langle F, \int_0^t d\tilde{a}_s^- \tilde{a}_s^- G \rangle \\
&= \langle F, \int_0^t \tilde{a}_s^- d\tilde{a}_s^- G \rangle + \langle F, \int_0^t d\tilde{a}_s^- \tilde{a}_s^- G \rangle, \quad F, G \in \mathcal{S},
\end{aligned}$$

hence (6). Concerning (7),

$$\begin{aligned}
\langle \tilde{a}_t^+ F, \tilde{a}_t^+ G \rangle &= \int_0^t \langle \tilde{a}_s^+ F d\tilde{N}_s, G \rangle - \int_0^t \langle \tilde{a}_s^+ F, \tilde{D}_s G \rangle ds \\
&\quad + \langle F, \int_0^t \tilde{a}_s^+ G d\tilde{N}_s \rangle - \int_0^t \langle \tilde{D}_s F, \tilde{a}_s^+ G \rangle ds + \langle N_t F, G \rangle \\
&= \langle \int_0^t d\tilde{a}_s^+ \tilde{a}_s^+ F, G \rangle + \langle \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ F, G \rangle - \int_0^t \langle \tilde{a}_s^+ F, \tilde{D}_s G \rangle ds \\
&\quad + \langle F, \int_0^t d\tilde{a}_s^+ \tilde{a}_s^+ G \rangle + \langle F, \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ G \rangle \\
&\quad - \int_0^t \langle \tilde{D}_s F, \tilde{a}_s^+ G \rangle ds + \langle N_t F, G \rangle \\
&= \langle \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ F, G \rangle + \langle F, \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ G \rangle + \langle N_t F, G \rangle \\
&= \langle \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ F, G \rangle + \langle \int_0^t \tilde{a}_s^- d\tilde{a}_s^+ F, G \rangle + \langle N_t F, G \rangle, \quad F, G \in \mathcal{S}.
\end{aligned}$$

□

The commutation relation stated in [1], [10] now reads

$$[\tilde{a}_t^-, \tilde{a}_t^+] = \tilde{a}_t^- \tilde{a}_t^+ - \tilde{a}_t^+ \tilde{a}_t^- = N_t + \int_0^t \tilde{a}_s^- d\tilde{a}_s^+ - \int_0^t d\tilde{a}_s^+ \tilde{a}_s^- + \int_0^t d\tilde{a}_s^- \tilde{a}_s^+ - \int_0^t \tilde{a}_s^+ d\tilde{a}_s^-.$$

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